## Comparison of the T<sub>1</sub> and D<sub>1</sub> Diagnostics: A New Definition for the Open-Shell D1 Diagnostic

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#### **Abstract**

It is shown that the  $T_I$  operator used in a previous study to define the open-shell  $D_I$  diagnostic is invalid, and leads to an arbitrary definition of the open-shell  $D_I$ diagnostic. A new definition is proposed that eliminates this ambiguity and approximately restores the mathematical relationship previously noted between the closed-shell  $D_I$  and  $T_I$  diagnostics. Statistical comparison of the  $T_I$  and  $D_I$  diagnostics shows a very high degree of correlation between them, although it is argued that both diagnostics used together can provide more information than either can separately.

#### Introduction

Recently, Leininger et al. [1] defined the open-shell version of the  $D_I$  diagnostic, the closed-shell version of which was defined by Janssen and Nielsen in 1998 [2]. The  $D_1$ diagnostic is based on the matrix 2-norm of the coupled-cluster t<sub>1</sub> amplitudes, and, like the  $T_I$  diagnostic [3-6] which is based on the Frobenius norm, it is designed to give an indication of the quality of results to be expected from a singles and doubles coupledcluster (CCSD) calculation. Janssen and Nielsen showed that there was a rigorous mathematical relationship between the  $D_I$  and  $T_I$  diagnostics for closed-shell theory, which is due to the fact that they are both based on the coupled-cluster t<sub>1</sub> amplitudes. For the open-shell version, however, no such relationship exists even though both the  $D_1$ and  $T_I$  open-shell diagnostics were defined for coupled-cluster wavefunctions based on restricted Hartree-Fock reference functions. There is a difference, however, since the open-shell  $T_I$  diagnostic was defined using an open-shell coupled-cluster wavefunction that is an  $S_x$  eigenfunction [6] based on symmetric spin-orbitals [7], whereas Leininger et al. defined a new  $T_I$  operator based on symmetric and antisymmetric combinations of annihilation and creation operators. This  $T_I$  operator was not actually used in calculating the wavefunction, rather the standard  $S_z$  spin-orbital equations were used, and the resulting  $t_1$  amplitudes transformed into those for the newly defined  $T_1$  operator. We show below, however, that the new  $T_I$  operator is not valid and that the relationship between the  $S_z$  open-shell  $t_1$  amplitudes and the new amplitudes is arbitrary. Hence, the definition of the open-shell  $D_l$  diagnostic is also arbitrary.

One of the purposes of this study is to eliminate the ambiguity in the open-shell  $D_I$  diagnostic by basing it on the  $S_x$  restricted open-shell coupled-cluster wavefunction developed in Ref. [6]. The mathematical relationship between the  $D_I$  and  $T_I$  diagnostics is now valid for both the open- and closed-shell versions. We provide test examples of

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Relationships between the coupled-cluster amplitudes in the  $S_x$  and the usual  $S_z$  basis were also given in Ref. [6]. For current purposes, a brief summary is given here. The  $T_1$  operator, spin rotated into an  $S_z$  basis, is given by Eq. (45) in Ref. [6] and is

$$T_{1} = f_{i}^{a} E_{ai} + v_{i}^{a} A_{ai} + \sqrt{2} f_{x}^{a} \left[ \left( a \alpha \right)^{\dagger} \left( x \alpha \right) \right] + \sqrt{2} f_{i}^{x} \left[ \left( x \beta \right)^{\dagger} \left( i \beta \right) \right] , \tag{3}$$

where  $\mathbf{f}$  and  $\mathbf{v}$  are the  $\mathbf{t}_1$  amplitudes in the  $S_x$  basis. The  $\mathbf{f}$  amplitudes correspond to orbital relaxation parameters whereas the  $\mathbf{v}$  amplitudes correspond to a type of double excitation in the  $S_x$  basis (see Ref. [6] for a more complete discussion). Equating Eq. (1) and Eq. (3) allows us to determine relationships between the  $\mathbf{f}$  and  $\mathbf{t}$  amplitudes:

$$f_i^a = \left(t_{i\alpha}^{a\alpha} + t_{i\beta}^{a\beta}\right)/2 \quad , \tag{4}$$

$$\frac{J_i - (t_{i\alpha} + t_{i\beta})^{1/2}}{f_r^a = t_{r\alpha}^{a\alpha} / \sqrt{2}}, \tag{5}$$

$$\begin{aligned}
f_x &= t_{x\alpha} / \sqrt{2} \\
f_i^x &= t_{i\beta}^{x\beta} / \sqrt{2}
\end{aligned} ,$$
(6)

$$\frac{J_i}{v_i^a} = \frac{t_{i\beta}}{t_{i\alpha}^{a\alpha} - t_{i\beta}^{a\beta}} / 2 \quad ,$$
(7)

where Eq. (7) is given for completeness. The open-shell  $D_I$  diagnostic is then defined as the matrix 2-norm of the  $\mathbf{f}$  amplitudes analogous to the procedure used in Ref. [1]. That is, the three open-shell amplitude blocks are treated separately and the maximum matrix 2-norm is taken as the  $D_I$  diagnostic in order to preserve the energy invariance properties of the open-shell  $D_I$  diagnostic (see Ref. [1]).

$$D_{1} = \max(\|f_{i}^{a}\|_{2}, \|f_{x}^{a}\|_{2}, \|f_{i}^{x}\|_{2}) .$$
(8)

Again, following Ref. [1], the matrix 2-norm is computed according to

$$\|\mathbf{R}\|_{2} = \sqrt{\lambda_{\text{max}}} = \sigma_{\text{max}} \quad , \tag{9}$$

where  $\lambda_{\max}$  is the largest eigenvalue of the matrix  $\mathbf{R}\mathbf{R}^T$  and  $\sigma_{\max}$  is the largest singular value of  $\mathbf{R}$ . Similar to the relationship between the open- and closed-shell  $T_I$  diagnostics, if there are no open-shell orbitals x, then the open-shell  $D_I$  diagnostic is identical to the closed-shell definition [2]. Defining the open-shell  $D_I$  diagnostic in this way has the additional advantage that the equation relating the closed-shell  $D_I$  and  $T_I$  diagnostics,  $\sqrt{2}T_1 \leq D_1$ , now approximately holds for the open-shell diagnostics in the limit of the number of open-shell electrons being much smaller than the total number of electrons. Test case examples of the new  $D_I$  diagnostic are given in the next section followed by a critical comparison to the  $T_I$  diagnostic.

### Results and Discussion

# A. New Open-Shell D<sub>1</sub> Values

Test case examples of the new open-shell  $D_I$  diagnostic together with the open-

both the closed-shell and open-shell molecules included in these two studies. To give an idea of how strong the correlation between  $T_1$  and  $D_1$  is, we can compare to the correlation coefficient between the  $T_1$  and  $S_2$  diagnostics ( $S_2$  is the perturbation theory equivalent of the  $T_1$  diagnostic [11]), as well as to the correlation coefficient between the coupled-cluster and perturbation theory  $D_1$  diagnostics [2]. Using the data for closed-shell molecules contained in Ref. [2], we obtain a correlation coefficient of 0.98 between the  $T_1$  and  $S_2$  values, and a correlation coefficient of 0.95 between the coupled-cluster and perturbation theory  $D_1$  diagnostics. This shows that the correlation between the  $D_1$  and  $T_1$  coupled-cluster diagnostics is almost as strong as the correlation with their respective second-order perturbation theory analogues.

Perhaps this high degree of correlation is to be expected because of the similarities between the two diagnostics. Both the  $T_l$  and  $D_l$  diagnostics are based on the orbital relaxation parts of the t<sub>1</sub> amplitudes, both were designed to exhibit the same orbital invariance properties that the CCSD energy possesses, and in spite of suggestions to the contrary [1,2], they were both designed to exhibit the mathematical property of sizeintensivity. For example, both diagnostics will yield the same value for a single helium atom as for any number of non-interacting helium atoms. However, it was pointed out in Ref. [5] that the  $T_I$  diagnostic may fail to indicate that a small region of a large molecule is difficult to describe properly if the rest of the molecule is well described at the CCSD level of theory, and in Ref. [2] Janssen and Nielsen gave a numerical example of this phenomenon. In essence, the  $T_I$  diagnostic is an average over the whole molecule, and the contribution from the small problem area is swamped by that from the majority of the molecule, which is well described. This can be viewed as a failure of the  $T_I$  diagnostic or it can also be viewed as a success since the majority of the molecule is well described and the  $T_I$  diagnostic indicates this. Conversely, the  $D_I$  diagnostic is designed to yield a large value for a large molecule with only a small problem area. Again, this could be viewed as a success or a failure since most of the molecule is well described, which the  $D_I$  diagnostic does not suggest, but there is one problem area which the  $D_I$  diagnostic does indicate. Another situation where the two diagnostics could give conflicting information is when there is an accumulation of correlation effects wherein there is no one orbital relaxation parameter or excited state that is very important, but rather there are several states or orbital relaxation parameters that are moderately important. In this case, the  $T_I$  diagnostic would be larger than usual, indicating the need to treat higher-order correlation, which would be missed by the  $D_I$  diagnostic. The important point, is that together the  $T_I$  and  $D_l$  diagnostics provide more information than either does alone, and it is best to use both. We stress, however, that we also believe that diagnostics based on two-particle components of the wavefunction, such as the  $D_2$  diagnostic [12], can provide additional, important information in assessing the quality of a particular calculation.

Another quantity to consider is the ratio  $T_l/D_l$ , and these are included in Table 1 for the open-shell molecules studied here. In a perfectly homogeneous system, this ratio can be derived from the approximate mathematical relationship between the  $T_l$  and  $D_l$  diagnostics to be  $1/\sqrt{2}$ . A homogeneous system in this sense is defined as one in which the contribution from all of the **f** amplitudes is identical. Thus, the further the  $T_l/D_l$  ratio deviates from  $1/\sqrt{2}$ , the greater the non-homogeneity of the electronic structure of the molecule being studied. We note that the  $T_l/D_l$  ratio will deviate from a perfectly

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